

Effective field theory of inflation confronts CMB and LSS: a progress report

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- ❑ **Effective field theory of inflation: condensate + fluctuations**
- ❑ **Small parameters: H/M_p + slow roll \longrightarrow expansion in $1/N_e$**
- ❑ **Fluctuations: initial conditions, backreaction and the Quadrupole**
- ❑ **Model analysis and reconstruction program.**

Effective field theory basics

A single scalar field

$$\mathcal{L} = \frac{1}{2} \partial_\mu \Phi \partial^\mu \Phi + V(\Phi)$$

$$\Phi(\vec{x}, t) = \phi + \delta\Phi(\vec{x}, t)$$

condensate

fluctuation

$$\delta\Psi \sim \frac{\delta T}{T}$$

$$g_{\mu\nu}(\vec{x}, t) = FRW + \delta g_{\mu\nu}(\vec{x}, t)$$

Tensor (grav.waves)

Vectors

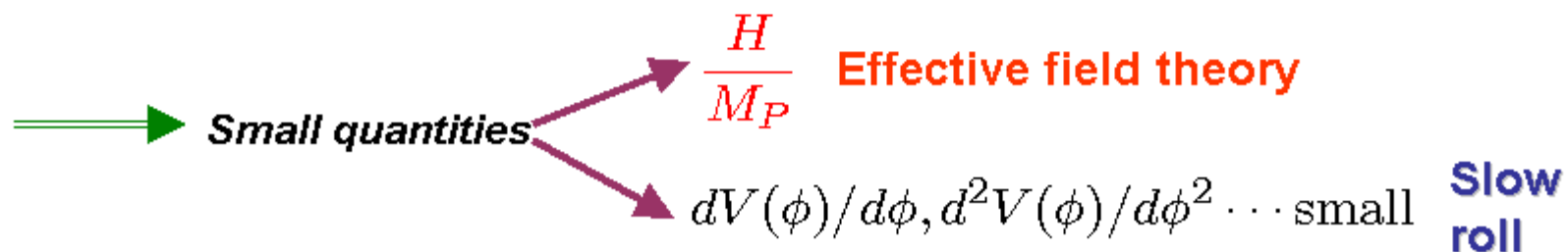
Scalar

$$H^2 = \frac{8\pi}{3M_P^2} \left[\frac{\dot{\phi}^2}{2} + V(\phi) \right]$$

$$\ddot{\phi} + 3H\dot{\phi} + V'(\phi) = 0$$

$$\delta\ddot{\Psi}_k + 3H\delta\dot{\Psi}_k + \frac{k^2}{a^2(t)}\delta\Psi_k + \mathcal{W}(\phi)\delta\Psi_k = 0$$

- Gravitational waves amplitude $\sim \frac{H}{M_P} \ll 10^{-5}$
- Nearly scale invariant spectrum of temperature fluctuations
- Need > 50 e-folds to solve horizon....problems



$$\epsilon_v = \frac{M_P^2}{2} \left[\frac{V'(\phi)}{V(\phi)} \right]^2, \eta_v = M_P^2 \frac{V''(\phi)}{V(\phi)}, \dots$$

Slow roll parameters $\ll 1$

Curvature perturbations (Gaussian):

Boundary conditions and Quadrupole suppression

COBE, WMAP1, WMAP3 \longrightarrow *low quadrupole*

$v_k = \delta\Psi_k a(t) \dot{\phi} / H$ obeys a wave eqn. in conformal time η

during slow roll

$$\left[\frac{d^2}{d\eta^2} + k^2 - \frac{\nu^2 - \frac{1}{4}}{\eta^2} \right] v_k = 0; \nu = \frac{3}{2} + 3\epsilon_v - \eta_v$$

Solutions with Bunch-Davies b.c.

$$g_\nu(k; \eta) = \frac{1}{2} i^{\nu+\frac{1}{2}} \sqrt{-\pi\eta} H_\nu^{(1)}(-k\eta); g_\nu(k; \eta) \stackrel{\eta \rightarrow -\infty}{=} \frac{1}{\sqrt{2k}} e^{-ik\eta}$$

General solution: $S(k; \eta) = A(k) g_\nu(k; \eta) + B(k) [g_\nu(k; \eta)]^*$

With normalization condition $|A(k)|^2 - |B(k)|^2 = 1$

Quantization:

$$v_k(\eta) = a_k S(k; \eta) + a_k^\dagger S^*(k; \eta) \quad B(k)=0 \longrightarrow \text{B.D. vacuum}$$

Power spectrum

$$\mathcal{P}(k) = \langle 0 | |\delta\Psi_k|^2 | 0 \rangle = \mathcal{P}_{BD}(k) [1 + D(k)]$$

$$\frac{H^2}{\epsilon_v M_P^2} \left(\frac{k}{k_0} \right)^{n_s - 1}$$

Transfer function for boundary conditions

$$n_s = 1 - 6\epsilon_v + 2\eta_v$$

Conditions on $D(k)$:

- Negligible backreaction on Einstein's eqns.
- Finite $T_{\mu\nu}$
- Renormalization of $T_{\mu\nu}$ with counterterms indep. of b.c

—————▶ **$D(k)$ falls faster than $1/k^4$**

Change in b.c. —————▶ change in C_l

$$\frac{\Delta C_l}{C_l} = \frac{\int_0^\infty D(\kappa x) f_l(x) dx}{\int_0^\infty f_l(x) dx}$$

$$\kappa = \frac{a_0 H_0}{3.3}$$

$$f_l(x) = x^{n_s - 2} [j_l(x)]^2$$

Rapid fall-off of $D(k)$ —————▶ **ONLY LOW MULTIPOLES ARE AFFECTED**

Q: What is the origin of $D(k)$?? (what determines the b.c.??)

A: a fast-roll stage PRIOR to slow roll

Allowing for **RAPID** variation of the condensate ϕ

$$\left[\frac{d^2}{d\eta^2} + k^2 - \frac{\nu^2 - \frac{1}{4}}{\eta^2} - \mathcal{V}(\eta) \right] v_k = 0$$

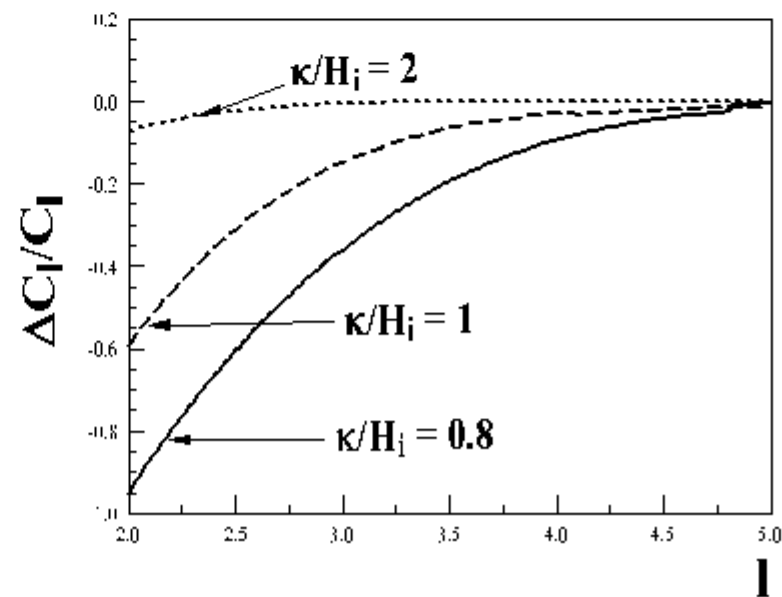
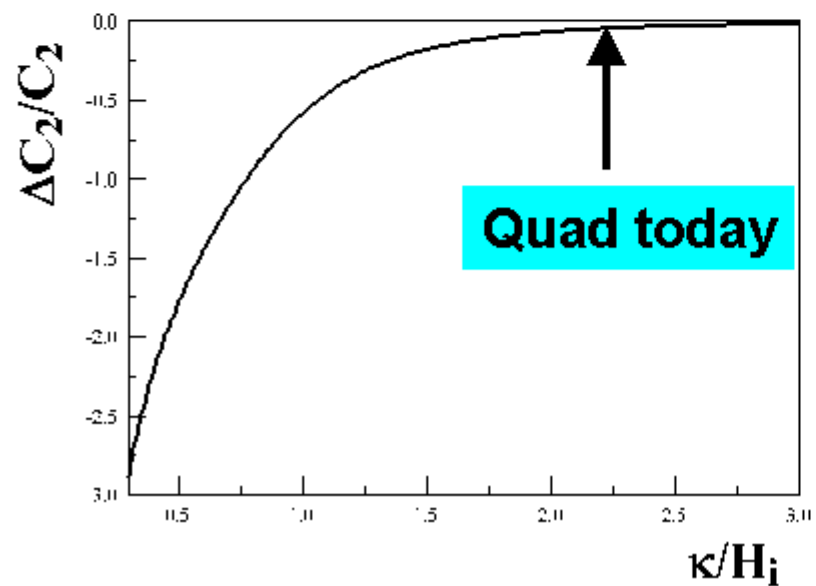
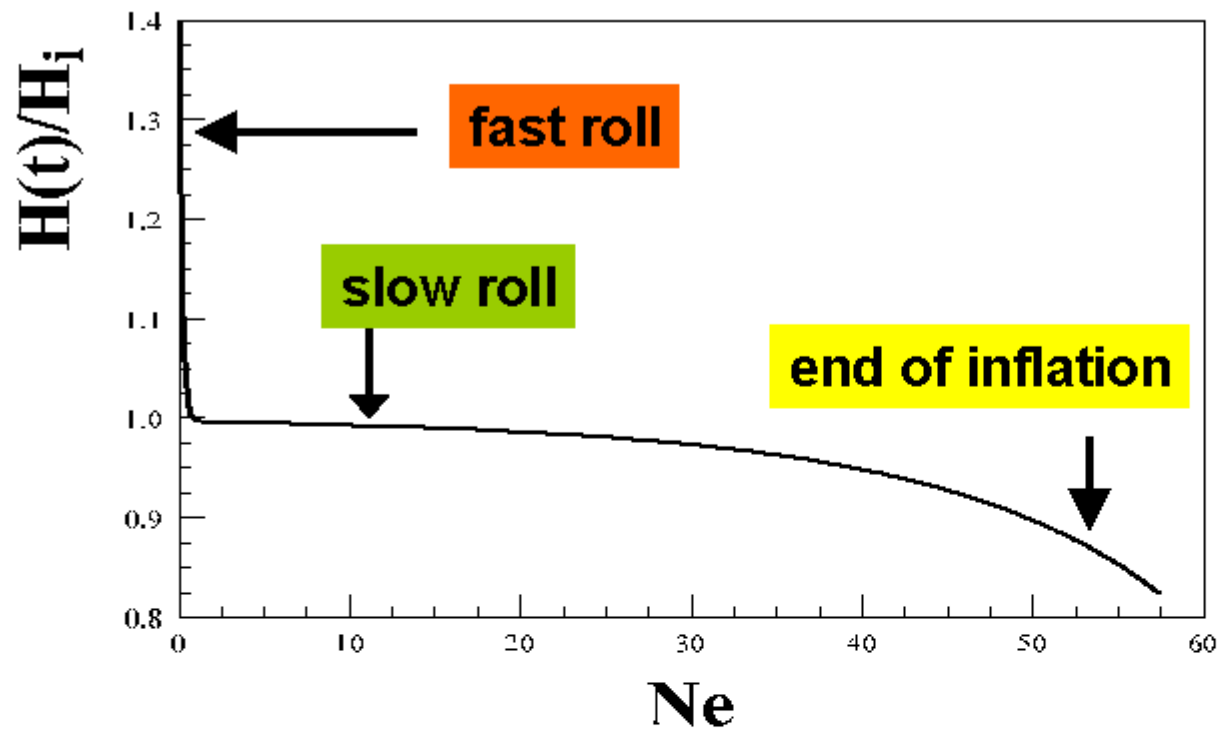


Depends on high(er) derivatives of ϕ : **negligible in slow roll**, **large for fast roll**

WHEN? large INITIAL $\dot{\phi}$ but large FRICTION term \rightarrow **short fast roll stage**

$\rightarrow \mathcal{V}(\eta) =$ **LOCALIZED POTENTIAL**

$D(k) \propto T(k)$ = Transmission coeff. of SCATTERING PBM!!



Executive summary:

- ❖ Initial conditions: set by a *fast roll* stage prior to slow roll
- ❖ Fast roll ~ **GENERIC** initial condition with kin. ~pot. energy
- ❖ Fast roll: localized *attractive* potential in mode equations
- ❖ $D(k)$ = transmission coeff. = transfer function
- ❖ Suppression of low multipoles

Analysis: fast roll ~ 2-3 e-folds, IF k_Q crosses horizon ~ 2-3 e-folds after beginning of slow roll ~ 15-20% suppression!

Quad. suppression mechanism *WITHIN* eff. field theory

**New vs. chaotic inflation and reconstruction program:
confronting *WMAP 3***

Implement eff. field theory + slow roll as $1/N_e$ expansion to systematically explore a large “family” of inflaton potentials.

WMAP 3 + LSS:

$n_s = 0.958 \pm 0.016$ (assuming $r = 0$ with no running)

$r < 0.28$ (95% *CL*) no running

$r < 0.67$ (95% *CL*) with running

Slow Roll expansion: a hierarchy of dimensionless parameters:

$$\epsilon_v = \frac{M_P^2}{2} \left[\frac{V'(\phi)}{V(\phi)} \right]^2, \quad \eta_v = M_P^2 \frac{V''(\phi)}{V(\phi)} \dots$$

as a $1/N_e$ expansion:

$$N[\phi(t)] = -\frac{1}{M_P^2} \int_{\phi(t)}^{\phi_{end}} V(\phi) \frac{d\phi}{dV} d\phi$$

$$\phi = \sqrt{N_e} M_P \chi \quad \longleftarrow \text{Rescale field}$$

$$V(\phi) = N_e M^4 w(\chi) \quad \sim \mathcal{O}(1)$$

energy scale of inflation

$$1 = - \int_{\chi_c}^{\chi_{end}} \frac{w(\chi)}{w'(\chi)} d\chi$$

value of χ at the END of inflation

value of χ N_e e-folds before END of inflation

Independent of N_e : an implicit relation between χ and couplings

$$\epsilon_v = \frac{1}{2 N_e} \left[\frac{w'(\chi_c)}{w(\chi_c)} \right]^2, \quad \eta_v = \frac{1}{N_e} \frac{w''(\chi_c)}{w(\chi_c)}$$

Explicit dependence on N_e : SIMPLE RESCALING

$$n_s - 1 = -6\epsilon_v + 2\eta_v \sim 1/N_e$$

$$r = 16\epsilon_v \sim 1/N_e$$

$$\frac{dn_s}{d \ln k} = -\frac{2}{N_e^2} \left\{ \frac{w'(\chi_c)w'''(\chi_c)}{w^2(\chi_c)} + 3 \left[\frac{w'(\chi_c)}{w(\chi_c)} \right]^4 - 4 \frac{[w'(\chi_c)]^2 w''(\chi_c)}{w^3(\chi_c)} \right\}$$

Simple scaling with $1/N_e$, choose $N = 50$ as representative

Family of potentials

$$V(\phi) = V_0 - \frac{1}{2} m^2 \phi^2 + \frac{\lambda}{2n} \phi^{2n}, \text{ broken symmetry}$$

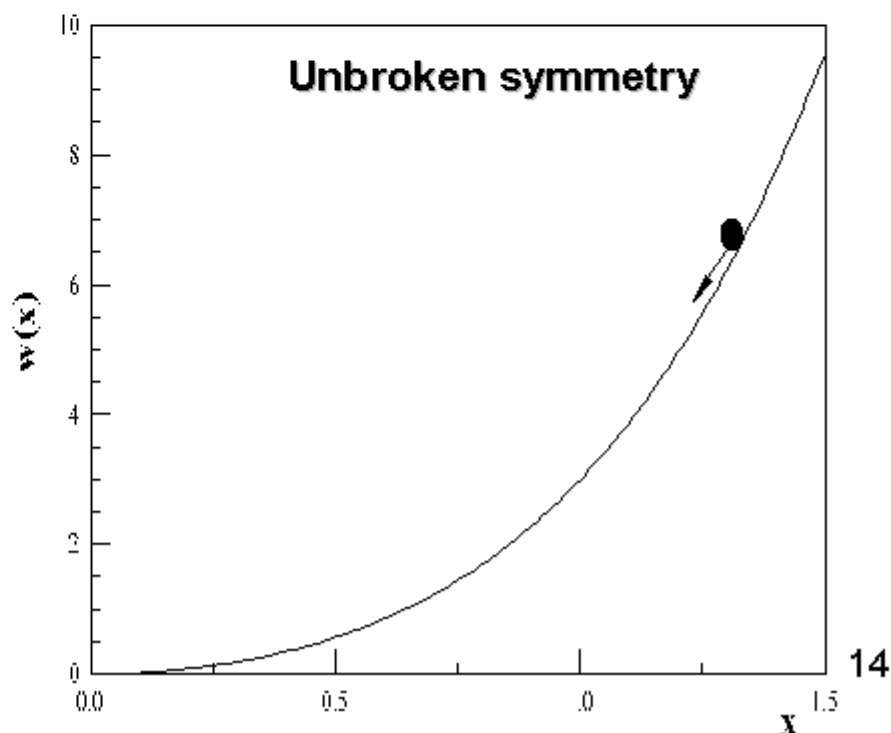
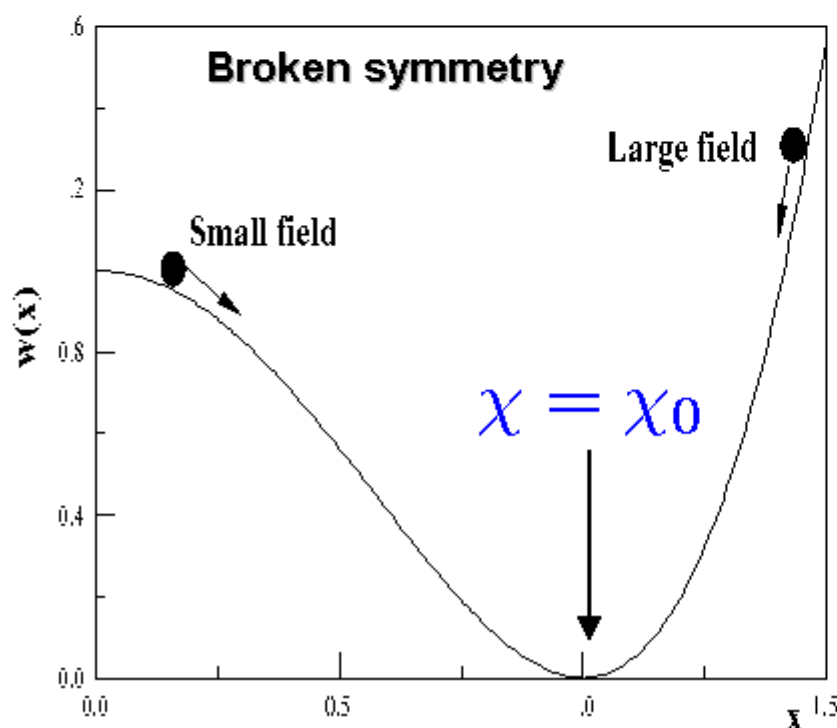
$$V(\phi) = \frac{1}{2} m^2 \phi^2 + \frac{\lambda}{2n} \phi^{2n}, \text{ unbroken symmetry}$$

Rescale fields and couplings:

$$\lambda = \frac{m^2 g}{M_{Pl}^{2n-2} N_e^{n-1}}; \quad g = \frac{1}{\chi_0^{2n-2}}; \quad x = \frac{\chi}{\chi_0}$$

$$w(\chi) = \frac{\chi_0^2}{2n} [n(1-x^2) + x^{2n} - 1]$$

$$w(\chi) = \frac{\chi_0^2}{2n} [n x^2 + x^{2n}]$$



Broken Symmetry

$$\frac{2n}{\chi_0^2} = \int_X^1 \frac{dx}{x} \frac{n(1-x^2) + x^{2n} - 1}{1-x^{2n-2}}$$

Unbroken symmetry

$$\frac{2n}{\chi_0^2} = \int_0^X \frac{n + x^{2n-2}}{1 + x^{2n-2}} x dx$$

$$X = \frac{\chi_c}{\chi_0}$$

Conditions for number of e-folds

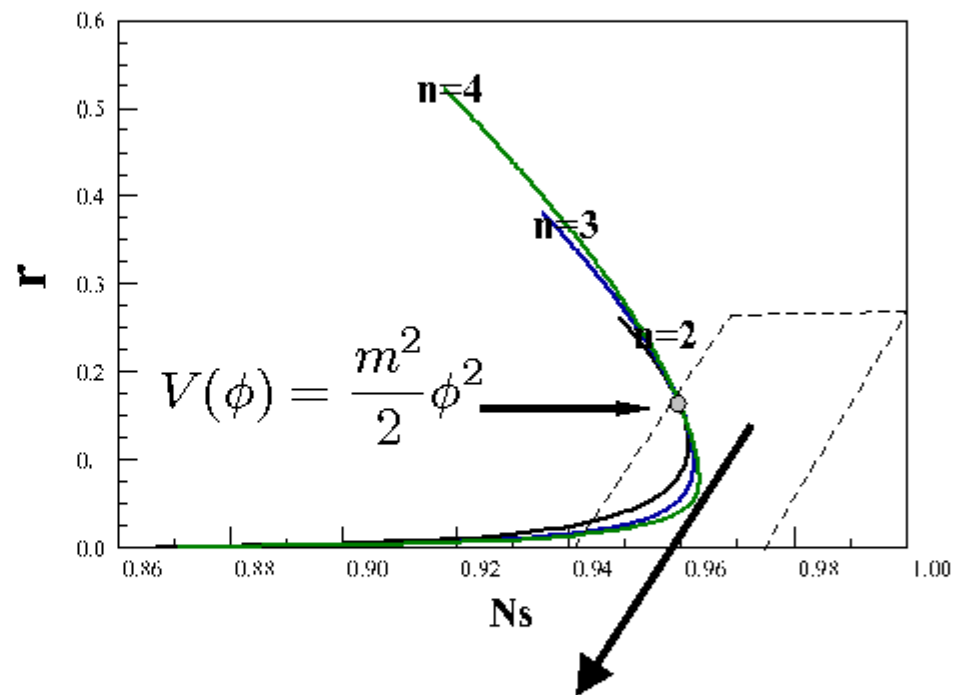
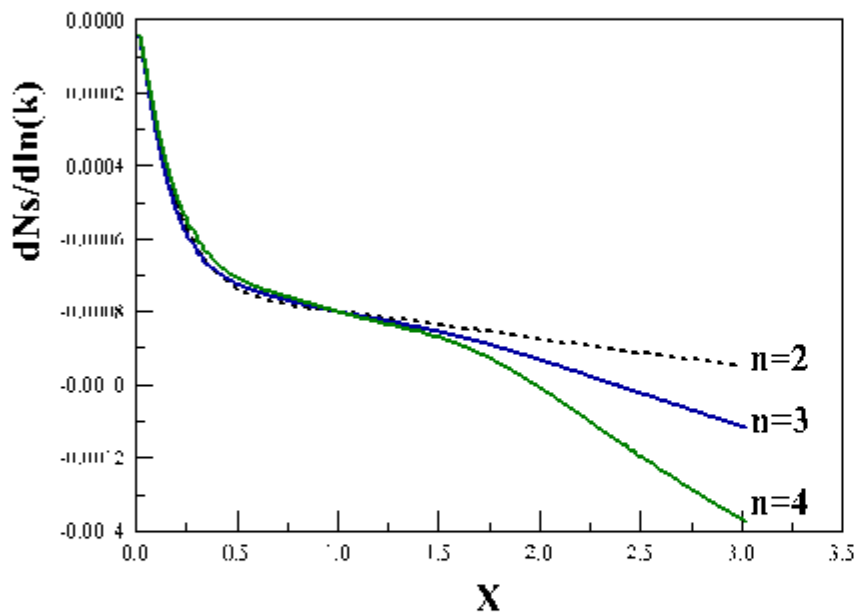
- X determines the value of the field N_e e-folds before the end of inflation
- χ_0 determines the value of the coupling for which the value of X is N_e before the end of inflation

STRATEGY

- 1) Vary X , find χ_0 construct $w[x]$ and derivatives
- 2) Find $\epsilon_v, \eta_v, n_s, r, dn_s/d \ln k$ as a function of X
- 3) Plot parametrically

RESULTS

1) New inflation (B.S.) $N_e = 50$ (change accordingly)

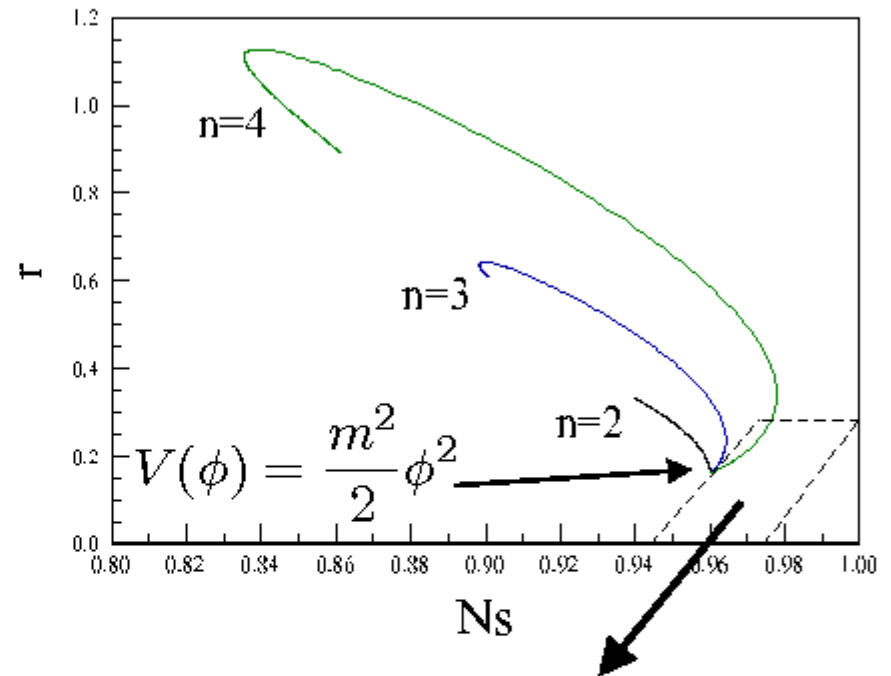
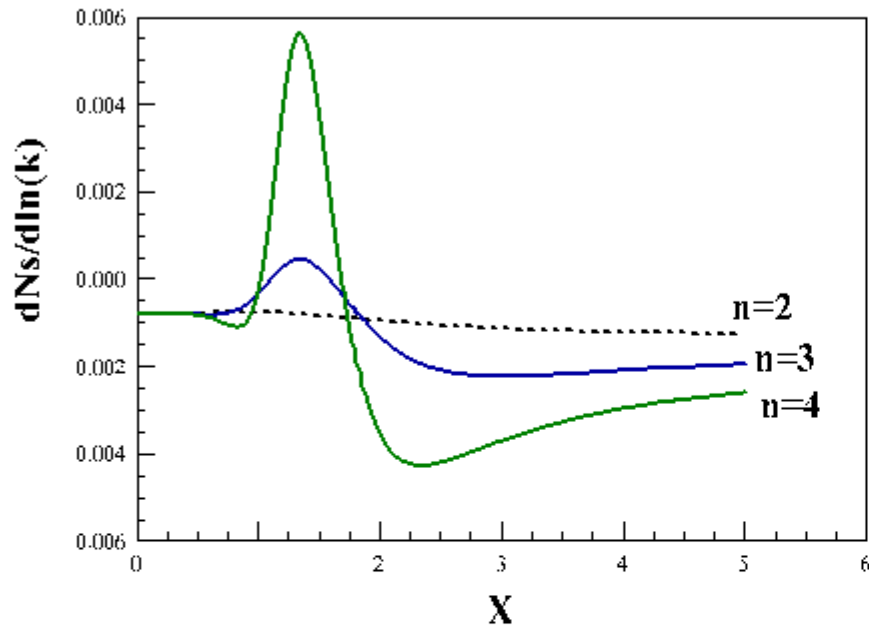


WMAP 3 marginalized region of r - n_s (95% CL)

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Large region of consistency for small field New Inflation

2) Chaotic inflation



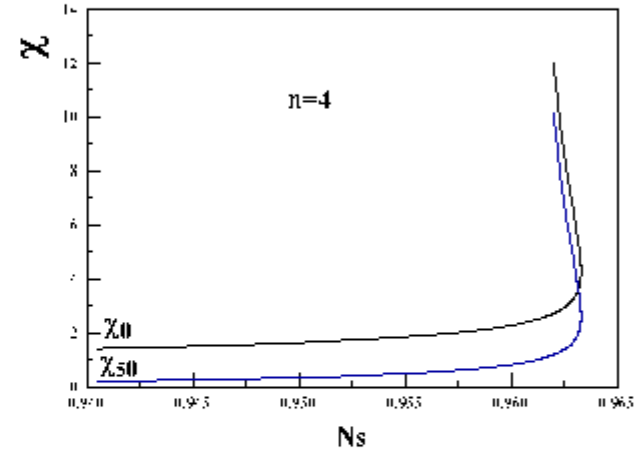
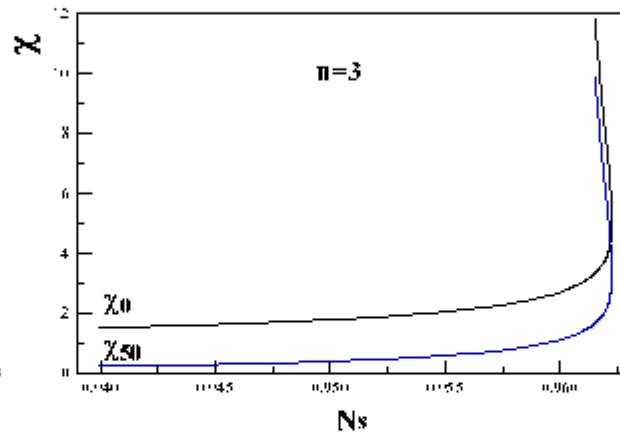
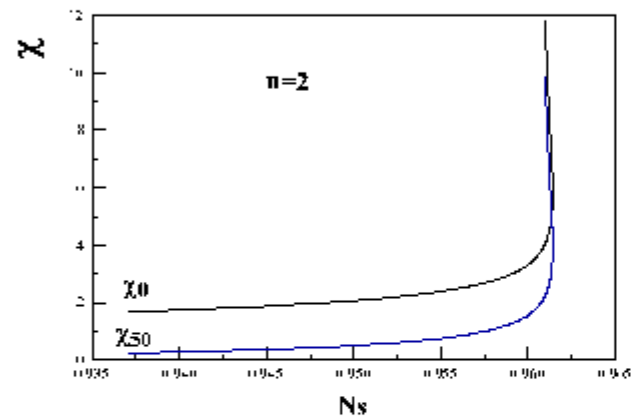
WMAP 3 marginalized region of r - n_s (95% CL)



Small region of consistency with WMAP 3

FIELD RECONSTRUCTION: $N_e = 50$

New inflation small field: region of consistency with WMAP 3



Symmetry breaking scale: $\phi_0 \sim 10 M_P$

Crossing scale: $\phi_{50} \sim M_P$

FINAL—FINAL SUMMARY AND CONCLUSIONS

- ✓ *Effective field theory $H/Mp \ll 1$, $1/N_e$ -slow roll expansion robust, systematic, predictive*
- ✓ *Quantum corrections suppressed by $(H/Mp)^2$*
- ✓ *Fast roll stage prior to slow roll \longrightarrow modifies b.c. for scalar perturbations \longrightarrow quadrupole suppression $\sim 15\text{-}20\%$ for total $N_e \sim 55$.*
- ✓ *$1/N_e$ expansion \longrightarrow systematic exploration of family of inflaton potentials+ field reconstruction .*
- ✓ *Small field New Inflation larger region of consistency with WMAP3+LSS data.*
- ✓ *Potentials with larger overlap with marginalized WMAP 3 data symmetry breaking scale $\sim 10 Mp$, crossing scale $\sim Mp$.*